

# Capillary Pressure Induced CO<sub>2</sub> Retention

Y. B. Altundas, R. de Loubens and T. S. Ramakrishnan

Schlumberger-Doll Research

Ridgefield, CT 06877

Geological sequestration involves injection of CO<sub>2</sub> into depleted oil and gas fields, saline aquifers, subsea sediments and deep coal beds. Among these, injection into saline aquifers appears to have the maximum storage potential. Various mechanisms that enhance long term storage include geological trapping, dissolution and mineralization.

In addition to the above mentioned four mechanisms, large volume of CO<sub>2</sub> may be immobilized, although the injected fluid may remain largely connected. This retention mechanism is driven by capillary pressure hysteresis. This is different from residual CO<sub>2</sub> trapping that arises through disconnections caused by fluid imbibition.

## 1 Introduction

A significant increase in atmospheric CO<sub>2</sub> concentration has been observed since the onset of the last century. The present level of CO<sub>2</sub> concentration in the atmosphere is reported to be 381 ppm, the highest ever seen in last 400 thousand years [1][2]. Several studies have shown that the change in CO<sub>2</sub> concentration and other greenhouse gases cause a warming effect necessitating technologies that mitigate CO<sub>2</sub> accumulation [1]. One of the technologies is geological carbon sequestration.

Geological sequestration involves injecting captured CO<sub>2</sub> directly into depleted oil-gas reservoirs, saline aquifers, and unminable coal beds. While saline aquifer sequestration has no tangible benefits, it has by far the largest storage potential. CO<sub>2</sub> injected into saline aquifers is expected to migrate slowly updip until containment by impermeable boundaries. In the absence of barriers, the time scale for migration should be kept sufficiently large to allow dissolution into saline water, eventually trapping CO<sub>2</sub> (permanently). Other trapping mechanisms that have been suggested are mineralization and residual CO<sub>2</sub> rich phase via counter imbibition [3][4].

In addition to the above, we propose a new migration retardation mechanism that we call hysteresis induced immobilization. Due to capillary pressure hysteresis, a gradient in  $\text{CO}_2$  saturation may be sustained essentially forever by having zero radial capillary pressure gradients and therefore possibly zero phase pressure gradients.

This paper illustrates hysteresis induced immobilization. A quantitative model of scanning curves is given. The model is sufficiently simple that relatively fast algorithms may be deployed. Computational techniques that are robust in handling capillary hysteresis are also incorporated. Naturally, the model includes relative permeability hysteresis as well. The negative permeability is completely consistent with the relative permeability model. Thus, in addition to geological trapping, our simulation includes two modes of sequestering  $\text{CO}_2$  residual phase formation, and immobilization via the hysteresis mechanism. We do not include the solubilization in this paper.

As an example, we construct a model problem to show hindered movement of a nonwetting phase upon cessation of injection. A comparison of results with and without hysteresis delineates the additional effect caused by the capillary pressure differences during retraction and injection. We also propose a quantitative measure for computing the effect of hysteresis on the movement of  $\text{CO}_2$ .

## 2 Mathematical Model

In this work, we emphasize the hysteresis effect on  $\text{CO}_2$  sequestration. We neglect some of the effects that are secondary to this issue, e.g., eventual solubility of  $\text{CO}_2$  in water. The reservoir is isothermal and is water-wet in relation to  $\text{CO}_2$ .  $\text{CO}_2$  is represented as a nonwetting phase with a constant compressibility, with the compressibility specified at the nominal reservoir pressure. More, importantly, we assume that  $\text{CO}_2$  is immiscible with water, but salt and water are perfectly miscible. We also adopt the work on zero volume of mixing in which it is assumed that there is no volume change of mixing between saturated brine and pure water. The description of this approach can be found in the appendix [5] [7].

For clarity, using subscripts, we denote the wetting phase (aqueous phase) by  $a$  and the  $\text{CO}_2$  phase by  $o$ . Thus, the mass balance equations for  $\text{CO}_2$  and aqueous phases can be given as follows.

$$\frac{\partial}{\partial t}[\phi \rho_o S_o] + \nabla \cdot (\rho_o \mathbf{v}_o) = 0 \quad (1)$$

$$\frac{\partial}{\partial t}[\phi \rho_a S_a] + \nabla \cdot \left[ \rho_a \mathbf{v}_a + \delta \frac{\mathcal{D} \nabla \psi}{1 + \delta \psi} \right] = 0 \quad (2)$$

where  $P$ ,  $\rho$  and  $S$  represent the corresponding pressure, density and saturation for CO<sub>2</sub> and aqueous phases, respectively. Also,  $\psi$  is volume fraction of the saturated salt solution in the brine,  $\mathcal{D}$  is the diffusion coefficient and  $\phi$  is the porosity,  $\mathbf{v}$  are the velocities.  $\delta$  is  $(\rho_a^s - \rho_a^o)/\rho_a^o$  where  $\rho_a^s$  is the saturated brine density, and  $\rho_a^o$  is the density of water, i.e., without salt. At the reservoir temperature and user-specified nominal reservoir pressure, CO<sub>2</sub> compressibility is specified using

$$c_o = \frac{1}{\rho_o} \left( \frac{\partial \rho_o}{\partial P_o} \right) \quad (3)$$

The superficial velocity  $\mathbf{v}_\beta$  is given by Darcy's law:

$$\mathbf{v}_\beta = -\frac{k k_{r\beta}}{\mu_\beta} (\nabla P_\beta - \mathbf{g} \rho_\beta), \quad \beta = a \text{ or } o \quad (4)$$

where  $k_{r\beta}$  and  $\mu_\beta$  are the relative permeability and the viscosity in  $\beta$ - phase, and  $\mathbf{g}$  and  $k$  are the acceleration due to gravity and the permeability of the formation, respectively. An explicit form of relative permeabilities and capillary curve functions by Ramakrishnan and Wasan[6] are implemented.

Based on the zero volume of mixing between saturated brine and water, the mass balance equation for salt can be given as follows

$$\frac{\partial(\phi\psi S_a)}{\partial t} + \nabla \cdot \left[ \psi \mathbf{v}_a - \frac{\mathcal{D} \nabla \psi}{1 + \delta \psi} \right] = 0 \quad (5)$$

With the inclusion of the following relations

$$P_c - P_o + P_a = 0 \quad (6)$$

$$S_o + S_a = 1 \quad (7)$$

we have a complete system of equations for 2-phase 3-component flow in porous media that we will use for CO<sub>2</sub> injection and monitoring.

## 2.1 Boundary conditions and wellbore implementation

The top and bottom boundary conditions of the reservoir are considered to be impermeable. Thus there is no flux through these boundaries, i.e.,

$$\mathbf{v}_\beta \cdot \mathbf{n}_\beta = 0, \quad \beta = a, o \quad (8)$$

which yields the following top and bottom boundary conditions for  $\psi$  as well, i.e.,

$$\frac{\partial \psi}{\partial z} = 0 \quad (9)$$

For the far-field boundary conditions, it is assumed that the outer boundary will be far from the inlet and

the pressures at the far-field will not be altered during the injection. Thus, we set  $P_\beta = P_{\beta_\infty}$ ,  $\beta = a, o$ , but no salt flux is allowed through the outer boundary. The inlet boundary condition for pressures and salt concentration are the injection or production flow rates top of the well.

In this work, wellbore is part of the reservoir but with orders of magnitude larger permeability than the formation layers. Characteristic wellbore capillary pressure is therefore very small, and in the presence of both phases, the wellbore will have near equality of phase pressures.

## 2.2 Initial conditions

Specifying a free water level, the densities as a function of pressure, and a capillary pressure relation

$$S_a = S_{ra} + (1 - S_{ra}) \left( \frac{P_b}{P_c} \right)^\lambda \quad (10)$$

we obtain the initial aqueous phase saturation. The salt concentration is let to be uniform.

## 3 Numerical Method

In order to discretize the equations (1)-(7), we use fully implicit block centered control volume method and the time derivative,  $\frac{\partial P_c}{\partial t}$ , is discretized via forward finite difference. Also, for computing the aqueous and CO<sub>2</sub> phase flow rates,  $Q_\beta$ , we use upstream weighting for relative permeabilities.

Fully implicit discretization is unconditionally stable but it leads to excessive numerical dispersion. With explicit discretization, dispersion is reduced, but one is forced to use smaller time steps to meet the CFL stability criterion. We, therefore, use adaptive implicit methods (AIM). The idea behind AIM is to discretize Eq. 5 explicitly at the grid points where time steps are acceptably large and meet the CFL condition [11]. Every where else, discretization will be implicit. In Fig. 2, we have the diagram showing the iterative process for solving the Eqs. 1-7.

## 4 Results

We compare our results with TOUGH2. As a reference we have considered a single salt-transport problem in a cylindrical geometry. In this problem, salt is injected into fresh water in a radial reservoir. The reservoir is 20 m thick and 33 m long. The injection rate is  $10^{-5}$  m<sup>3</sup>/s and stopped after 100 days. The initial salt concentration in the reservoir is about 17%(wt). In Fig. 3, we show the contour plots obtained from TOUGH2 and our code (cFAST). The results are very comparable and, moreover, cFAST shows less numerical dispersion in relation to TOUGH2 for the same number of grids.

## 4.1 Effect of capillary pressure

The discussion below emphasizes the effect of capillary pressure on the movement of a nonwetting phase such as supercritical CO<sub>2</sub>. During the injection of CO<sub>2</sub> into saline aquifer,  $S_o$  will increase as high as  $1 - S_{ra}$  near the injection well. The CO<sub>2</sub> plume migrates upward due to buoyancy. After the cessation of injection, the upward movement is compensated by counter imbibition. Locations where the salination reversal occurs are impeded by having to follow capillary pressure appropriate to that location, i.e., the prior history determines the  $P_c$  and  $k_r$  functions for each location. In our model, the prior history is captured by keeping track of the lowest  $S_a$  reached for that location. Recent studies by Orr *et al.* [12] have taken hysteresis into account, although their numerical computations are centered on relative permeability hysteresis as opposed to ours where both relative permeability and capillary pressure loops are considered. The capillary pressure hysteresis dominates the induced immobilization, once residual saturation are accounted for.

After cessation of CO<sub>2</sub> injection, CO<sub>2</sub> plume will move upward due to lower density, and saline water will imbibe (see Fig. 4). CO<sub>2</sub> migrating upward is a drainage process whereas, saline water replacing CO<sub>2</sub> is an imbibition process. Here, the capillary pressure curve is reversed from drainage to imbibition curve, thus different parts of the formation follow different legs of the scanning curves.

In this work we use the relative permeability and capillary pressure formulas that take into account the disconnection of nonwetting phase, and integrates both capillary pressure and relative permeability hysteresis in a simple way (see appendix). The underlying relative permeability functions are the same as those of Land [13]. The capillary pressure hysteresis is introduced via two mechanisms: trapping effects and pore body to pore throat ratio ( $\alpha$ ) [6].

## 4.2 Test Problem

In order to simulate  $P_c$  hysteresis induced trapping, we consider a cylindrical reservoir where outer radius is at 2000 m with a thickness of 30 m. The reservoir properties are given in Table. 1. The reservoir has a single injection well and is homogeneous.

A slice of the domain is shown in Fig. 5. CO<sub>2</sub> injection is carried out at a rate 0.17 m<sup>3</sup>/sec for one year. At the end of first year, the injection is stopped. The computation was however was carried out for a total of 50 years. After injection, CO<sub>2</sub> migrates upward due to lower density and CO<sub>2</sub> is replaced through counter-imbibition of brine. Due to negligible viscous pressure drops, during this phase of storage, details of capillary pressure hysteresis becomes dominant. The simulation is carried out primarily to study

the influence of hysteresis over time scale relevant to storage.

To distinguish the  $P_c$  hysteresis trapping from others, we use a very quantitative approach in which we prescribe a region to monitor the change in  $\text{CO}_2$  saturation and calculate the amount of  $\text{CO}_2$  remaining in the prescribed region. To choose such region, we first calculate the radius,  $r_c$ , of the porous cylinder with a given height and porosity so that cylinder's pore volume is equal to the volume of injected  $\text{CO}_2$ . The precise choice of this radius is not important. In this work, the radius of such a cylinder is about 543 m but we set it to be 545 m. Any movement of  $\text{CO}_2$  tongue reduces the amount of  $\text{CO}_2$  within this cylinder.

By varying the strength of the capillary pressure hysteresis through  $\alpha$ , and keeping everything else constant, The enhanced  $\text{CO}_2$  immobilization may be quantified. Changing  $\alpha$  alters the magnitude of the area within the drainage and imbibition capillary pressure curves, without affecting relative permeabilities. Note that  $\alpha$  being the ratio of pore body to pore throat size, the influence of the pore geometry is directly reflected in the computations.

Table. 2 shows the amount of  $\text{CO}_2$  present within a radius of 545 m at the end of 50 years. The retention volume of  $\text{CO}_2$  monotonically increases with  $\alpha$ . The difference between these choices in terms of volumetric capacity may be 30,000 tons around the well, not an insignificant amount. The evolution of this immobilized volume for various values of  $\alpha$  is shown in Fig. 6. Beyond a value of  $\alpha = 3$ , there appears to be little sensitivity to  $\alpha$ . Figures. 8 and 7 are the 3-D plots of aqueous saturations and  $\text{CO}_2$  saturations around the wellbore. Here,  $z = 0$  is the top of the wellbore. It is evident that the aqueous phase saturation far away from the wellbore has decreased when  $\alpha = 1$  due to a gravity tongue of  $\text{CO}_2$ . Figure. 7 clearly demonstrates this at the leading front of  $\text{CO}_2$ .

## 5 Summary

We have developed a numerical simulation that reduces numerical dispersion by using AIM, while taking into account: 1. Wellbore interaction and 2. Capillary pressure and relative permeability hysteresis.

The numerical simulation clearly shows that in addition to models of trapping such as leaking faults and residual phase formation, hysteresis induced retardation of flow could be a significant factor. This has strong implications when one is concerned about  $\text{CO}_2$  migrating updip and leaking through an outcrop. Any retardation of the  $\text{CO}_2$  tongue, increases chances of permanent storage by allowing additional time for solubilization.

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## 6 Appendix

In our algorithm, layer permeabilities and porosities are specified. As a default the entry capillary pressure  $P_b$  is related to an  $O(1)$  number  $J_b$  through

$$\frac{P_b}{\gamma} \sqrt{\frac{k}{\phi}} = J_b (\simeq 0.2) \quad (11)$$

where  $\gamma$  is the interfacial tension.

The drainage capillary pressure curve is given by the Brooks-Corey relation [14]

$$S_a^* = \left( \frac{P_b}{P_c} \right)^\lambda ; \forall P_c > P_b \quad (12)$$

where  $\lambda$  is the pore size distribution index and  $S_a^*$  is the normalized wetting phase saturation equal to  $(S_a - S_{ra})/(1 - S_{ra})$ . For all  $P_c \leq P_b$ ,  $Sw^* = 1$ .

We now define  $S_{wc} = (1 - S_{oc})$  as the lowest saturation reached during nonwetting phase intrusion. Then with

$$\begin{aligned} S_{ro}^* &= \frac{S_{ra}}{(1 - S_{ra})} \\ C &= \frac{1 - S_{rmo}^*}{S_{rmo}^*} \\ S_{wc}^* &= \frac{S_{wc}}{1 - S_{ra}} \end{aligned} \quad (13)$$

and the Land's relation that

$$S_{ro}^* = \frac{S_{oc}^*}{1 + C S_{oc}^*} \quad (14)$$

the disconnected phase saturation  $S^{*dc}$  of any stage of imbibition has been shown to be [6]

$$\begin{aligned} S_o^{dc*} &= \left[ C^{-1} S_{oi}^* + (1 - S_a^*)(C^{-1} + S_{oi}^*) - \right. \\ &\quad \left. \sqrt{[C^{-1} S_{oi}^* + (1 - S_a^*)(C^{-1} + S_{oi}^*)]^2 - 4C^{-2}(C^{-1} + S_{oi}^*)(S_{oi}^* + S_a^* - 1)} \right] / 2(C^{-1} + S_{oi}^*) \\ S^* &= S_a^* + S_o^{dc*} \end{aligned} \quad (15)$$

The above references also describe relative permeability as a function of  $S_w$ . The following are for drainage and imbibition relative permeability curves.



## A Nonwetting phase relative permeabilities ( $\lambda = 2$ )

The relative permeability of the nonwetting phase during *drainage* cycle is given by

$$k_{ro} = (1 - S_a^{*2})(1 - S_a^{*2}) \quad (16)$$

During *imbibition* cycle it is given by

$$k_{ro} = S_o^{c*2}(1 - S^{*2}) \quad (17)$$

where  $S_o^{c*} = S_o^* - S_o^{dc*}$

## B Wetting phase relative permeabilities ( $\lambda = 2$ )

The relative permeability of the wetting phase during *drainage* cycle is given by

$$k_{ra} = S_a^{*4} \quad (18)$$

The relative permeability of the wetting phase during *imbibition* cycle is given by

$$k_{ra} = 2S_a^{*2} \left[ \frac{S^{*2}}{2} - C^{-2} \left( \text{Log} \left( \frac{1 + C^{-1} - S^*}{C^{-1} + S_{oi}^*} \right) + (1 + C^{-1}) \frac{(S_{oi}^* + S^* - 1)}{(1 + C^{-1} - S^*)(C^{-1} + S_{oi}^*)} \right) \right] \quad (19)$$

## C Capillary pressure curve

In addition to trapping induced hysteresis represented by relative permeabilities, an algorithm for  $P_c$  maybe given as follows. If  $r_i$  is the pore body radius and  $r_d$  is the pore throat radius (drainage) then we replace the Brooks-Corey relation

$$S_a^* = \left( \frac{P_b}{P_c} \right)^{1/\lambda} \quad (20)$$

with a more general relationship

$$S^* = \left( \frac{P_b}{P_c \nu} \right)^{1/\lambda} \quad (21)$$

where  $S^* = S_a^* + S_o^{dc*}$ .

$S_o^{dc*} = 0$  in the initial drainage cycle. But during imbibition, the addition of  $S_o^{dc*}$ , allows us to capture

the correct pore entry radius corresponding to a given saturation  $S_a$  because the disconnected nonwetting phase is surrounded by the wetting phase. But in addition to that we have to take into account that the entry criterion for imbibition is dictated by the pore body whereas it is governed by pore throat for drainage. Therefore  $\nu = \alpha$  for imbibition and  $\nu = 1$  for drainage.

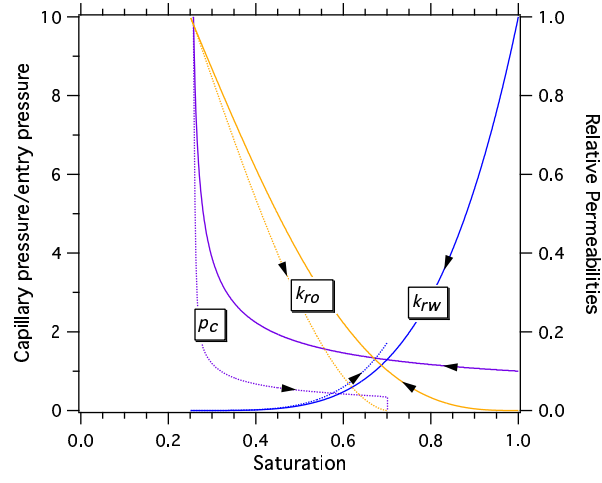


Figure 1: Capillary pressure and relative permeability curves

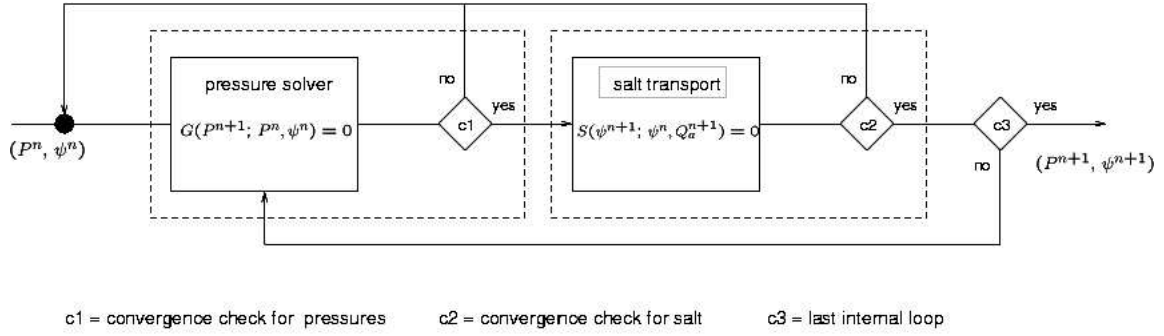


Figure 2: The flow diagram of the iterative process

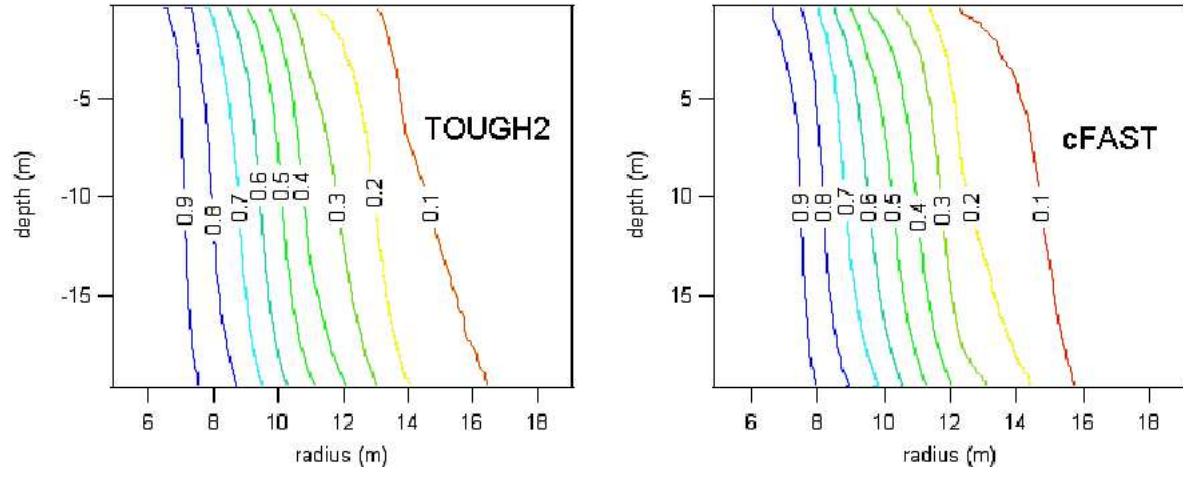


Figure 3: Figures show the contour plot of  $\psi$  obtained from TOUGH2 and cFAST.

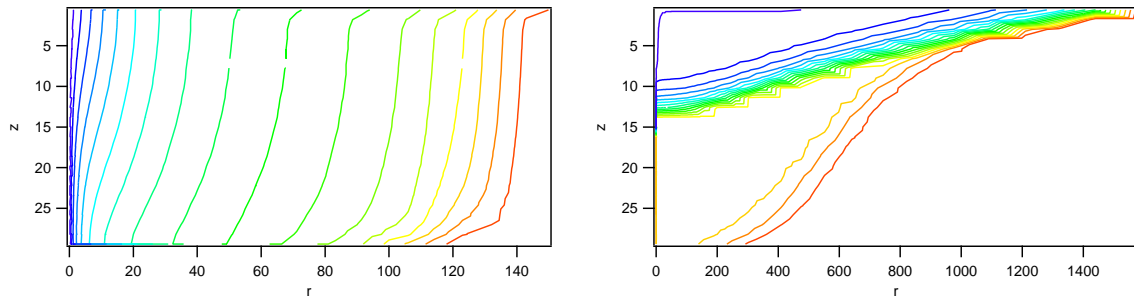


Figure 4: Counter plots of CO<sub>2</sub> saturation near wellbore in 10 days and 50 years

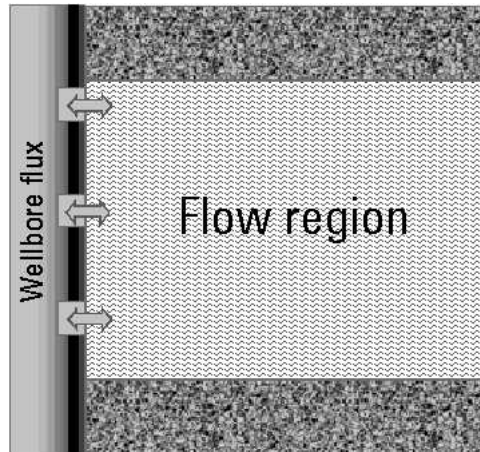


Figure 5: Wellbore and formation

Table 1: Reservoir properties

$k_v(\text{vertical})$	$10^{-13} \text{ m}^2$
$k_v(\text{horizontal})$	$10^{-13} \text{ m}^2$
$\phi$	20%
$q_{\text{inj}}$	$0.17 \text{ m}^3 \text{ CO}_2/\text{sec}$
$S_{ra}$	0.1
$S_{ro}$	0.2
$\psi_{\text{init}}$	0.34
$P_{\text{reference}}$	20 MPa
$T_{\text{reservoir}}$	60°C
$\mu_a(\text{water})$	0.47 cp
$\mu_a(\text{saturated})$	0.91 cp
$\mu_o$	0.06 cp
$\rho_a(\text{water})$	$983 \text{ kg/m}^3$
$\rho_a(\text{saturated})$	$1175 \text{ kg/m}^3$
$\rho_o$	$724 \text{ kg/m}^3$
$\text{CO}_2 \text{ compressibility}$	$2.2\text{e-}8 (\text{N/m}^2)^{-1}$

Table 2: Volume and percent of  $\text{CO}_2$  remained in the reservoir within a radius of 545 m from the wellbore

$\alpha$	Volume of $\text{CO}_2$ ( $\text{m}^3$ )	% of $\text{CO}_2$
1	1,783,010	36.31
2	1,815,907	36.98
3	1,827,052	37.20
4	1,831,239	37.29
5	1,833,690	37.38

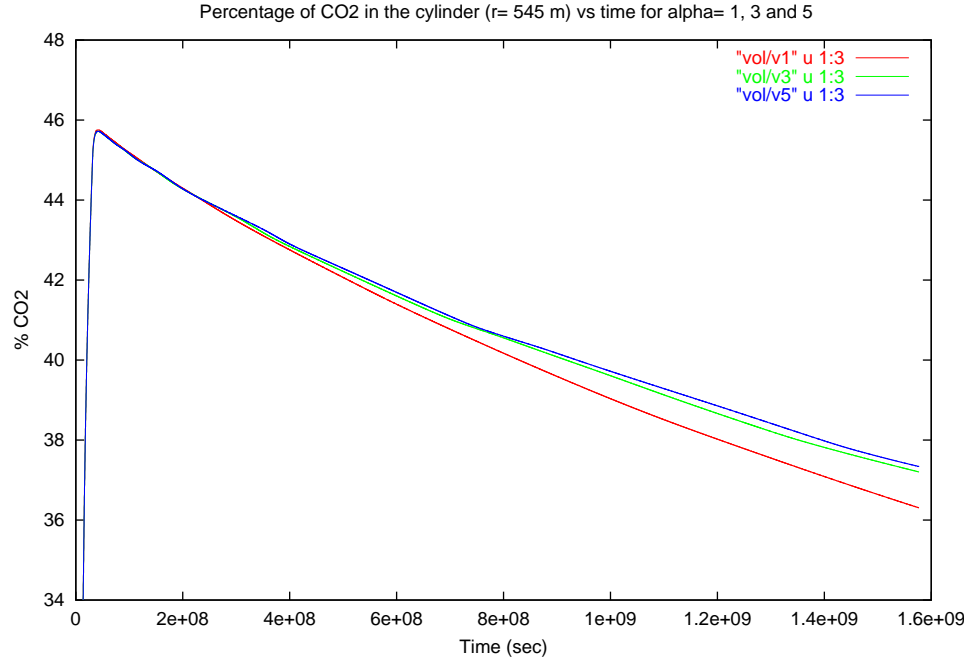


Figure 6: Figure shows the percentage of  $\text{CO}_2$  present in the cylinder vs the time for  $\alpha = 1, 2, 3, 4$  and  $5$

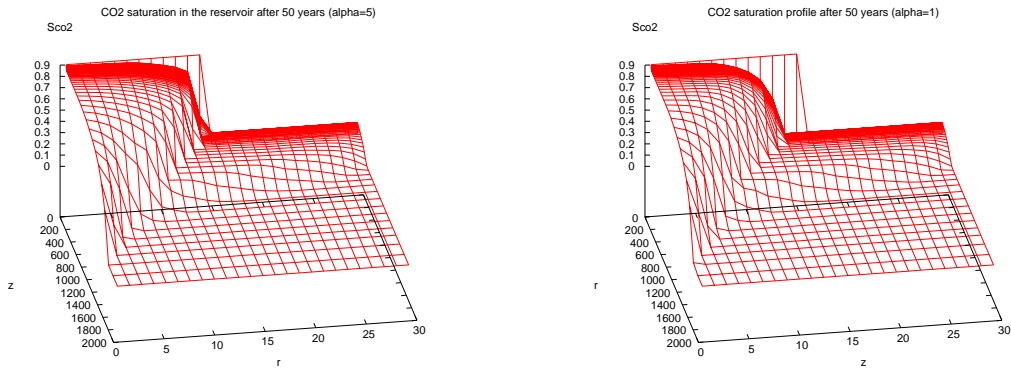


Figure 7:  $\text{CO}_2$  phase saturations 49 years after injection for  $\alpha = 1, 5$

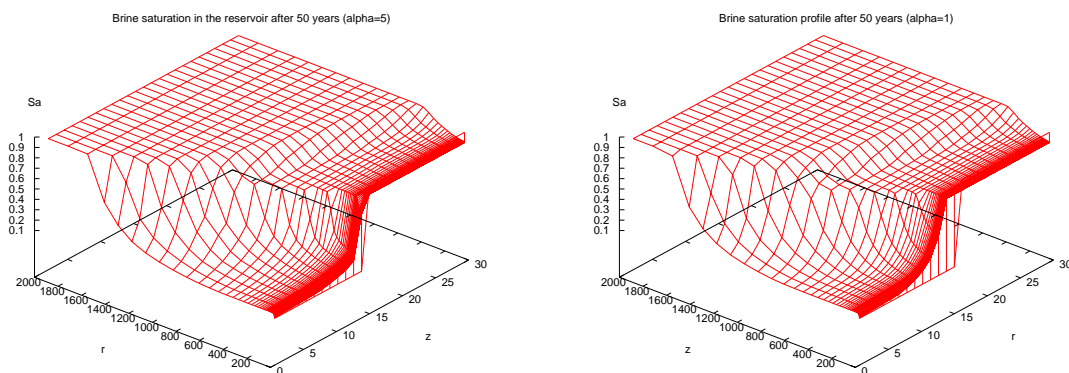


Figure 8: Aqueous phase saturations 49 years after injection for  $\alpha = 1, 5$

**T. S. Ramakrishnan**

Scientific Advisor

Schlumberger-Doll Research

36, Old Quarry Road

Ridgefield, CT 06877.

Ph: 1-(203)431-5239

Fx: 1-(203)438-3819

Em: ramakrishnan@slb.com

**R. de. Loubens**

Graduate student

Department of Petroleum Engineering

Stanford University

Em: rlouben@tobeadded

**Y. B. Altundas**

Senior Research Scientist

Schlumberger-Doll Research

36, Old Quarry Road

Ph: 1-(203)431-5527

Fx: 1-(203)438-3819

Em: baltundas@slb.com